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# Magneto-Optical Activity in Nonmagnetic Hyperbolic Nanoparticles

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Active nanophotonics can be realized by controlling the optical properties of materials with external magnetic fields. Here, we explore the influence of optical anisotropy on the magneto-optical activity in nonmagnetic hyperbolic nanoparticles. We demonstrate that the magneto-optical response is driven by the hyperbolic dispersion via the coupling of metallic-induced electric and dielectric-induced magnetic dipolar optical modes with static magnetic fields. Magnetic circular dichroism experiments confirm the theoretical predictions and reveal tunable magneto-optical activity across the visible and near infrared spectral range.

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Optical nanostructures enabling active control of light polarization are key components for future nanophotonic applications. One relevant example is the class of magneto-photon nanomaterials which exploit magneto-optical (MO) enhancement produced by either pure metallic or pure dielectric materials [1]. Artificial anisotropic metamaterials combining metals and dielectrics enable more versatile light-matter interaction engineering at the nanoscale [2,3], and their applications have diversified toward sensing [4,5], nanoscale waveguiding [6,7], enhanced nonlinearities [8–12], and amplified spontaneous emission [13–16]. The bulk optical properties of these materials can be described by an effective dielectric tensor, where the in-plane and out-of-plane permittivities,  $\epsilon_{\parallel}$  and  $\epsilon_{\perp}$ , respectively, can be almost arbitrarily tuned by the geometric dimension and material composition [17]. By stacking multiple layers of metal and dielectric materials, extreme optical anisotropies are

achieved, often described by hyperbolic isofrequency surfaces ( $\epsilon_{\parallel}\epsilon_{\perp} < 0$ ), which are opposed to the elliptical ones observed in typical materials like pure metals or dielectrics. This optical anisotropy can be further exploited to enhance the MO response of a material [18–21], allowing to control the optical properties by applying an external magnetic field, which is fundamental in view of the development of active magnetophotonic devices [22–27], from sensing [28–32], and all-optical switching [33,34], to nonreciprocal light propagation [35–38] and polarization selection [39,40].

Here, we explore the fundamental coupling of an external magnetic field to the localized electric and magnetic modes excited in nonmagnetic type-II ( $\epsilon_{\parallel} < 0, \epsilon_{\perp} > 0$ ) hyperbolic nanoparticles, where full control of the optical properties has been recently demonstrated [41,42]. The broadband MO response of the sample is studied by magnetic circular dichroism (MCD) spectroscopy. We observe direct modulation of the MO response at both fundamental resonances of the system, which we, based on numerical modeling, interpret in terms of magnetic field-induced spatial confinement or broadening of circular currents in the nanoparticles. By developing an analytical framework based on quasi-isotropic Mie theory, we show

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that the MO response can be fundamentally described on the single particle level by considering the coupling of dipolar resonances to a static magnetic field.

Disk-shaped hyperbolic nanoparticles with a nominal diameter of 250 nm were fabricated on fused silica substrates using hole-mask colloidal lithography [43] (see Supplemental Material [44]). Two types of hybrid nanoparticles were prepared, in which the dielectric material is varied between  $\text{TiO}_2$  and  $\text{SiO}_2$  to allow for tuning of the fundamental absorption resonance [41]. Five alternating layers of Au (10 nm each) and of the dielectric material (20 nm each) were deposited. A SEM image of the Au/ $\text{TiO}_2$  sample is shown in the inset of Fig. 1(a), alongside with the experimentally obtained optical extinction. Two peaks are clearly visible in the extinction spectrum, at  $\lambda \sim 1000$  nm and  $\lambda \sim 1600$  nm. Since the nanoparticles are randomly distributed, we do not expect any collective effect and can consider the measured extinction as the average optical response of a single nanoparticle. The unusual optical properties of these nanoparticles can be described by an effective medium theory, where the anisotropy is represented by an in-plane permittivity  $\epsilon_{\parallel}$  and an out-of-plane permittivity  $\epsilon_{\perp}$ . These can be expressed as [17]

$$\epsilon_{\parallel} = \frac{t_m \epsilon_m + t_d \epsilon_d}{t_m + t_d}, \quad (1)$$

$$\epsilon_{\perp} = \frac{\epsilon_m \epsilon_d (t_m + t_d)}{t_m \epsilon_d + t_d \epsilon_m}, \quad (2)$$

where  $t_{m(d)}$  is the thickness and  $\epsilon_{m(d)}$  is the permittivity of the metal (dielectric). The dispersion of the permittivity is shown as a function of the wavelength in the inset of Fig. 1(b). For  $\lambda > 500$  nm,  $\epsilon_{\perp}$  is positive (dielectriclike response) while  $\epsilon_{\parallel}$  is negative (metal-like response), corresponding to a type-II hyperbolic dispersion relation. To validate our approach,

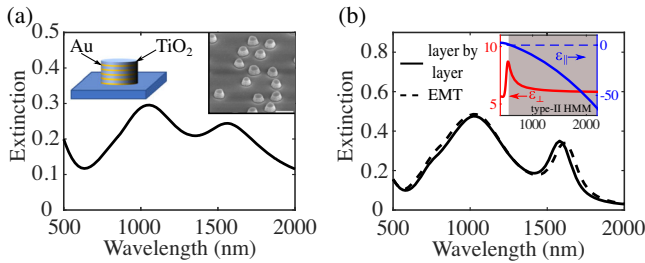


FIG. 1. (a) Experimental extinction spectrum of the Au/ $\text{TiO}_2$  hyperbolic nanoparticle metasurface. Insets are a sketch of the geometry (top left) and a SEM image of the sample (top right, scale bar is 400 nm). (b) Numerically obtained extinction spectra calculated based on the real structure (layer by layer, solid line) and using an effective medium theory (EMT, dashed line). Inset shows the real parts of the in-plane (blue) and out-of-plane (red) permittivities used in the EMT approach. For  $\lambda > 500$  nm, the system shows type-II hyperbolic dispersion (type-II HMM), as indicated in the figure (gray shaded area).

full-wave numerical simulations have been performed using tabulated permittivity values for gold [45] (see also Supplemental Material [44]). First, the optical extinction was calculated considering the real structure with alternating layers [solid line in Fig. 1(b)]. Afterward, we considered a homogenized structure described by an anisotropic hyperbolic effective permittivity as expressed in Eqs. (1) and (2) [dashed line in Fig. 1(b)]. The resulting spectra are in excellent agreement and describe well the experimentally measured extinction, demonstrating that the physical origin of these modes lays on the hyperbolic nature of the system. It is worth noticing that the high energy plasmon decays mainly radiatively, while the low-energy plasmon decays mainly non-radiatively, as can be seen in Supplemental Material, Fig. 1, thanks to the ability of these architectures to separately tune both decay channels, in contrast with conventional plasmonic architectures [cf. plain gold nanodisks, Supplemental Material, Fig. 2(a)].

When a magnetic field is applied, the moving charges undergo a Lorentz force  $\mathbf{F}_L = q(\mathbf{v} \times \mathbf{B})$ , oriented perpendicularly to the magnetic field and the trajectory of the charges. In this case, when a static field  $H$  is applied along, for instance, the  $z$  direction (out of plane), the symmetry of the dielectric function is broken, leading to nonzero off-diagonal components of the permittivity  $\hat{\epsilon}$  [50], i.e.,

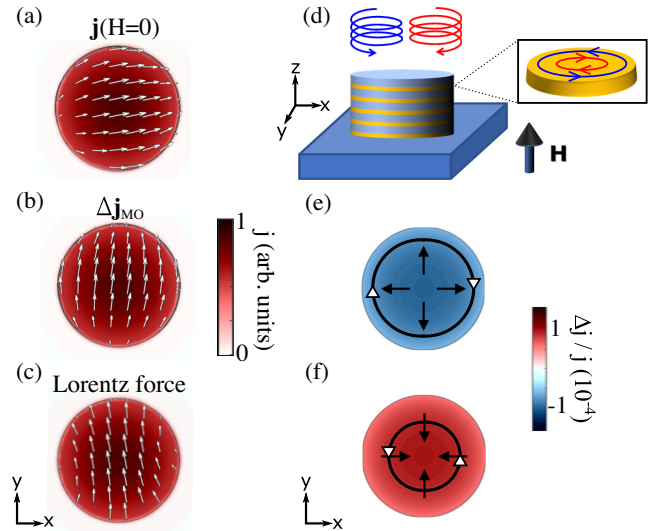


FIG. 2. (a)–(c) Calculated current density  $\mathbf{j}$  (white arrows) at the low-energy resonance for zero magnetic field (a), magnetically induced change of current  $\Delta \mathbf{j}_{\text{MO}}$  (b), and Lorentz force calculated as cross product of  $\mathbf{j}$  and the unit vector in direction of  $\mathbf{H}$  (c). The absolute value of  $j$  is color coded for all three cases. (d) Circular dichroism effect in terms of spatial confinement or broadening of circular modes. (e), (f) Induced relative change of current density  $\Delta j/j$  (color coded) for opposite helicity of the incoming electric field, as sketched in the plot. Black arrows show the gradient of  $\Delta j/j$  and display the induced spatial broadening or confinement of the circular plasmonic motion.

$$\hat{\epsilon} = \begin{pmatrix} \epsilon_{xx} & \epsilon_{MO} & 0 \\ -\epsilon_{MO} & \epsilon_{yy} & 0 \\ 0 & 0 & \epsilon_{zz} \end{pmatrix}. \quad (3)$$

The nondiagonal MO term  $\epsilon_{MO}$  is generally very small for noble metals compared to ferromagnetic materials (we assume  $\epsilon_{MO} = 10^{-4} + i10^{-3}$ , valid for Au at 1 T, whereas  $\epsilon_{MO} = 1 + i0.1$  for Co [50]). However, nanoscale design of bulk materials can drastically alter the MO properties due to the presence of plasmon resonances, leading to an enhanced MO response [50–52]. The MO coupling is particularly interesting in the context of circularly polarized light, where dichroism is induced by the applied magnetic field in systems that otherwise exhibit rotational symmetry in the propagation direction of the incoming light [26]. When a static magnetic field is applied along the propagation direction, the degeneracy of left- and right-handed modes is lifted due to the off-diagonal coupling term in Eq. (3), which is observable by a derivativelike spectral shape of the MCD signal. In conventional plasmonic nanomaterials, this magnetically induced mode splitting is observed at the fundamental electric dipole resonance (as shown for plain gold nanodisks in Supplemental Material, Fig. 2) [26,27]. Here, by considering the coupling of the charges with the magnetic field, we predict MCD also at the low-energy magnetic dipole mode, which is related to the dielectric nature of the system, hence rendering this resonance magnetoactive and thereby enriching the versatility of the anisotropic system in view of active functionalities. We numerically assess the induced changes in the current density inside a single nanoparticle due to the interaction with an external magnetic field. All quantities are evaluated inside the topmost gold layer of the structure, which is resonantly excited at the low-energy magnetic dipole mode. For this case, the numerically calculated current density  $\mathbf{j}$  at magnetic field  $H = 0$  is shown as white arrows in Fig. 2(a) for one phase of the incoming electric field together with the absolute value of the current density  $j$  (color coded). Because of the MO coupling in the gold layers, when a static magnetic field is applied in positive or negative  $z$  direction ( $+H$  and  $-H$ , respectively), we expect a magnetic-field dependent change in the current,  $\Delta\mathbf{j}_{MO} = \frac{1}{2}[\mathbf{j}(+H) - \mathbf{j}(-H)]$  [Fig. 2(b)]. Further, we directly compute the Lorentz force acting on the charges as the cross product of the current density and the magnetic field [Fig. 2(c)]. Indeed, the change in current density  $\Delta\mathbf{j}_{MO}$  follows well the direction of this force, as can be seen by comparing Figs. 2(b) and 2(c). One can notice orientation deviations of the arrows at the boundary of the particle, which is expected from the shape anisotropy of the system at its boundaries. An inversion of the magnetic field direction is topologically equivalent to a change in the helicity of the incoming light. Thus, for a fixed magnetic field and if averaged over one period of the electric field

oscillation, following the sketch presented in Fig. 2(d), the magnetically induced force on the free charge carriers will lead to spatial confinement or broadening of the circular modes, strictly depending on the helicity of the incoming electric field. This is verified via the change in absolute (i.e., independent on the phase  $\varphi$ ) current density  $\Delta j$  in Figs. 2(e) and 2(f), and consequently results in a shift of the resonance frequency for opposite helicities (i.e., circular dichroism).

To experimentally probe the interaction of the circular modes with an external magnetic field, we use MCD spectroscopy [46,47] (see Supplemental Material [44]). We define the MCD via the difference between two extinction spectra acquired using light with opposite helicity,  $E^+$  (left-handed circular polarization) and  $E^-$  (right-handed circular polarization) respectively, in the presence of an external magnetic field parallel to the incident light direction,

$$\text{MCD} = \frac{E^+ - E^-}{E^+ + E^-}. \quad (4)$$

Corresponding experimental MCD spectra for the Au/TiO<sub>2</sub> (red diamonds) and Au/SiO<sub>2</sub> (yellow circles) samples are shown in Fig. 3(a). A derivativelike spectral shape is clearly observed at the low and the high energy resonances for the Au/TiO<sub>2</sub> sample. Such line shape is consistent with the excitation of circular plasmonic modes in noble metal nanostructures with rotational symmetry in the plane of the incident electric field, as we show for plain gold nanodisks in Supplemental Material, Fig. 2. For the hyperbolic nanoparticles, the same line shape of the MCD spectrum is observed for both the electric and magnetic dipole modes, in agreement with the rationalization reported above, i.e., magnetically induced spatial confinement or broadening of circular currents in the nanoparticles upon circularly polarized light excitation (Fig. 2). Interchanging the dielectric layer with SiO<sub>2</sub> leads

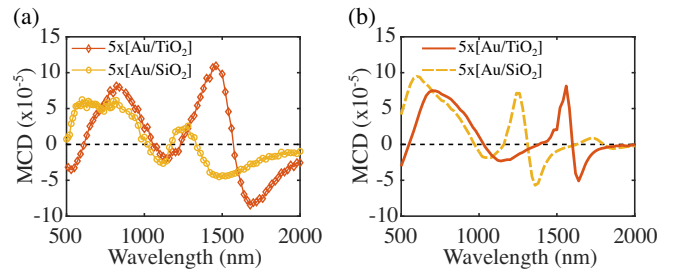


FIG. 3. (a) Experimental MCD spectra for the hyperbolic nanoparticles made of the two Au/dielectric multilayers: 5x[Au/TiO<sub>2</sub>] (red diamonds) and 5x[Au/SiO<sub>2</sub>] (yellow circles) samples acquired at 1.4 T applied magnetic field. (b) Simulated MCD spectra using the extended effective medium approach, including a MO coupling term, of the two Au/dielectric multilayers: 5x[Au/TiO<sub>2</sub>] (continuous line) and 5x[Au/SiO<sub>2</sub>] (dashed line).



to a strong blueshift of the low-energy resonance, shown by the yellow circles in Fig. 3(a) (see also Supplemental Material, Fig. 1).

To describe the MO properties of our system in the framework of an effective medium theory, Eqs. (1) and (2) are generalized as

$$\hat{\epsilon}_{\text{EMT}} = \begin{pmatrix} \epsilon_{\parallel} & f * \epsilon_{\text{MO}} & 0 \\ -f * \epsilon_{\text{MO}} & \epsilon_{\parallel} & 0 \\ 0 & 0 & \epsilon_{\perp} \end{pmatrix}, \quad (5)$$

where the off-diagonal component due to the coupling with the external magnetic field is considered for gold, and  $f = 1/3$  is the volume fraction of gold in the nanoparticle. Separate simulations are performed for both helicities and the MCD is calculated based on the individual extinction spectra according to Eq. (4). Resulting spectra are shown in Fig. 3(b) and reproduce well the observed spectral features. In these calculations, we assume a constant value of the MO coupling term,  $\epsilon_{\text{MO}} = 10^{-4} + i10^{-3}$ . Sepúlveda *et al.* further suggested a MO coupling term for the noble metals that is linear in the diagonal permittivity of the material,  $\epsilon_{\text{MO}} \propto \epsilon_{xx}$  [50]. This frequency dependent  $\epsilon_{\text{MO}}(\omega)$  again reproduces the expected derivativelike spectral shape, however it overestimates the dichroism effect at the low-energy resonance (Supplemental Material, Fig. 3).

The broadband MO response of the system fundamentally originates from the anisotropy of the dielectric function. The effect of this anisotropy is clearly observable in the numerically calculated electric near field of the resonantly excited nanoparticles (Supplemental Material, Fig. 4). At the high energy resonance, the major contribution originates from the in-plane electric field [Supplemental Material, Figs. 4(a) and 4(c)], while the out-of-plane contribution dominates for the low-energy resonance [Supplemental Material, Figs. 4(b) and 4(d)]. These results strongly suggest a dominant role of  $\epsilon_{\perp}$  and  $\epsilon_{\parallel}$  at the low and high energy resonances, respectively. To generalize this concept, we consider a spherical particle with radius  $r$  displaying either isotropic permittivity  $\epsilon_{\perp}$  or  $\epsilon_{\parallel}$  and calculate the scattering ( $\sigma_{\text{sca}}$ ) and absorption ( $\sigma_{\text{abs}}$ ) cross sections based on Mie theory [53,54]. The cross sections are calculated using a customized MATLAB code [48] as a weighted sum of the so-called Mie coefficients  $a_n$  and  $b_n$ , describing dipolar ( $n = 1$ ), quadrupolar ( $n = 2$ ), etc., vector spherical harmonics of electric ( $a_n$ ) and magnetic ( $b_n$ ) nature [49]. The calculated Mie efficiencies  $Q_{\text{abs/sca}} = \sigma_{\text{abs/sca}} / \pi r^2$  up to  $n = 1$  for  $r = 150$  nm are presented in Figs. 4(a) and 4(b) for the in-plane and out-of-plane components of  $\epsilon$ , respectively (for more details, see Supplemental Material [44]). To validate our approach, we compare the quasi-isotropic Mie theory with numerical results for a sphere with anisotropic permittivity  $\hat{\epsilon} = (\epsilon_{\parallel}, 0, 0; 0, \epsilon_{\parallel}, 0; 0, 0, \epsilon_{\perp})$  (Supplemental Material, Fig. 5).

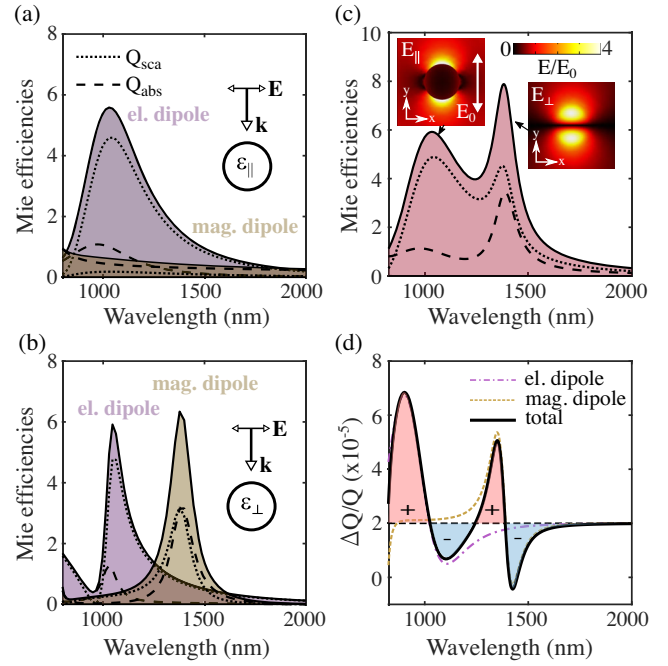


FIG. 4. Scattering (dotted lines) and absorption (dashed lines) efficiencies  $Q_{\text{sca}}$  and  $Q_{\text{abs}}$  for spherical nanoparticles with isotropic permittivity  $\epsilon_{\parallel}$  (a) and  $\epsilon_{\perp}$  (b). Interaction efficiencies are analytically calculated with the Mie expansion up to the first order, where magnetic (green) and electric (purple) contributions are separately evaluated. (c) Sum of the in-plane electric dipole shown in (a) and the out-of-plane magnetic dipole shown in (b). Inset are near-field plots of the resonances normalized to the incoming field  $E_0$  indicated by the white arrow. (d) Differential extinction spectrum  $\Delta Q/Q$  calculated from the magnetically induced dichroism effect. Contributions from the electric and magnetic dipoles are marked by the dash-dotted and dashed lines, respectively. The overall spectrum (black solid line) exhibits a derivative spectral shape at the resonances, as indicated by the light blue and light red areas under the curve.

As for the hyperbolic disks, the numerical extinction shows high and low-energy resonances which we can now allocate to an in-plane electric and an out-of-plane magnetic dipole, respectively. In our quasi-isotropic model, higher absorption at the low-energy resonance can be accounted for by adding a small fraction of the imaginary part of  $\epsilon_m$  to Eq. (2) (see parametric study in Supplemental Material, Fig. 6). The out-of-plane electric dipole is omitted here, because the in-plane permittivity contributes dominantly at that wavelength (Supplemental Material, Fig. 4). The sum of in-plane electric dipole and out-of-plane magnetic dipole is shown in Fig. 4(c) and agrees well with the numerical results (Supplemental Material, Fig. 5).

For noble metals like gold, the MO coupling is predominantly determined by free carriers in a broad spectral range [50] and the optical properties of the metal conduction electrons are well described by the Drude model. We consider the coupling of these charges to an external magnetic field in a perturbative approach [55], i.e., we

introduce a magnetic-field-dependent dielectric function of the metal,  $\epsilon_m = \epsilon_m^0 \pm \epsilon_{\text{MCD}}(H)$ . Since a change in the helicity of the incoming light is topologically equivalent to an inversion of the magnetic field direction, we can also understand the MCD in terms of a magnetically induced change of the plasmonic resonance condition for a fixed polarization of the incoming light. In our perturbative approach, we can express their MO coupling as [23]

$$\epsilon_{\text{MCD}}(H) = \frac{-2\omega_c(H)\omega_p^2\tau^3}{[\omega_c(H)^2\tau^2 - \tau^2\omega^2 + 1]^2 + 4\tau^2\omega^2}, \quad (6)$$

where  $\tau$  is the electron relaxation time,  $\omega_p^2 = Ne^2/\epsilon_0 m^*$  the plasma frequency, and  $\omega_c(H) = eH/m^*$  the magnetic-field-dependent cyclotron frequency ( $e$  is the electron charge,  $m^*$  the effective mass, and  $N$  the electron density). Inverting the applied magnetic field will thus change the orbital frequency of the free carriers, leading to a  $H$ -dependent modification of the plasmonic resonance condition, resulting in a spectral shift of the mode for an applied magnetic field. We can then express the induced MCD via the change in extinction caused by inversion of the magnetic field direction as

$$\frac{\Delta Q}{Q} = \frac{Q^{(+H)} - Q^{(-H)}}{Q^{(H=0)}}, \quad (7)$$

where  $Q^{(H=0)}$  is the extinction efficiency at  $H = 0$ . The corresponding spectrum is shown in Fig. 4(d) (black solid line). From the Mie expansion of the extinction, we can distinguish the contributions to the dichroism coming from either the electric (dash-dotted line) or the magnetic (dashed line) dipole resonance, thus obtaining a derivative-like spectral shape of the differential spectrum at both resonances, which describes well the induced dichroism observed in our experiments.

Finally, it is worth noticing that this system can be directly used also for active control of the polarization state of light via the polar MO Kerr effect without losing on modulation strength. Additionally, we can expect MO coupling in the longitudinal configuration, where the response should be one order of magnitude less strong. However, in the latter case nontrivial coupling between the different resonances might provide additional enhancement [35]. As well, we expect that all these MO effects can be enhanced by at least two orders of magnitude if the metallic building block is a magnetic material. Moreover, if the metallic layer is made of Co/Pt or Au/Co multilayers, this will help in reducing the applied field by almost one order of magnitude due to the strong magnetic anisotropy promoting magnetization along the out-of-plane direction [56,57]. Alternatively, the use of magnetic garnets as dielectric building blocks can also boost the MO response of these nanoparticles, as was previously demonstrated in iron garnets doped with gold nanoparticles and multilayered

magnetoplasmonic crystals [58–60]. Although the quality factor of the resonances reported in this work is not high and will further decrease when magnetic materials are employed, it was recently demonstrated for pure plasmonic nanoparticles that this factor can be improved, even by orders of magnitude compared to that of the constituent building blocks, by ordering the particles in a crystal-like structure and playing with their size and shape [61]. This approach will add interesting perspectives on how to combine coherent interparticle coupling, such as for surface lattice resonances, with anisotropic optical dispersion.

In conclusion, we experimentally and theoretically demonstrated MO activity in type-II hyperbolic nanoparticles by reasonably low ( $\sim 1$  T) external magnetic fields. Numerical simulations support the experimental results and prove the viability of the effective medium theory for the description of the anisotropic permittivity of these nanoparticles. Moreover, numerical modeling allows addressing the MO coupling in terms of induced changes in the current density of the system. We established a Mie-theory-based analytical model, calculating the MO response by considering the interaction of a static magnetic field with hyperbolic dispersion-induced magnetic and electric dipolar modes in individual nanoparticles. Beyond fundamental insights on the physics of the system, this analytical approach is general and can be expanded to other geometries and materials. Our results open exciting opportunities for the development of novel nanophotonic active materials combining the benefits of subwavelength light manipulation with the ability to freely tune the MO response via control over the anisotropy in the permittivity of the system.

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